


In the name of God

Engineering Mathematics

(Lecture # 10)

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Partial Differential Equations (PDEs)

Heat Equation: Very Long Bars

Heat Equation

(1)
$$\frac{\partial u}{\partial t} = c^2 \frac{\partial^2 u}{\partial x^2}$$

The results in the last section can be extended to bars of infinite length, which are good models of very long bars or wires (such as a wire of length, say, 100 m). Then the role of Fourier series in the solution process will be taken by **Fourier integrals**.

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Partial Differential Equations (PDEs)

(1) $\frac{\partial u}{\partial t} = c^2 \frac{\partial^2 u}{\partial x^2}$ → $u(x, t) = F(x)G(t)$

(2) $u(x, 0) = f(x) \quad (-\infty < x < \infty)$

$u(x, t; p) = [A(p) \cos px + B(p) \sin px] e^{-c^2 p^2 t}$

Using Fourier Integral.
 And Initial condition $u(x, 0) = f(x)$

(9) $u(x, t) = \frac{1}{\pi} \int_{-\infty}^{\infty} f(v) \left[\int_0^{\infty} e^{-c^2 p^2 t} \cos(px - pv) dp \right] dv.$

Remember these Results (9)&(10)!

(11) $u(x, t) = \frac{1}{2c\sqrt{\pi t}} \int_{-\infty}^{\infty} f(v) \exp\left\{-\frac{(x-v)^2}{4c^2 t}\right\} dv.$

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Use of Fourier Transforms

The **Fourier transform** is closely related to the **Fourier integral**.

$f(x) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \hat{f}(w) e^{iwx} dw$

$\hat{f}(w) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(x) e^{-iwx} dx.$

we can use Fourier transforms for solving our present or similar problems. The **Fourier transform** applies to problems **concerning the entire axis**, and the **Fourier cosine** and **sine transforms** to problems involving the **positive half-axis**.

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EXAMPLE 2 Temperature in the Infinite Bar in Example 1

Solve Example 1 using the Fourier transform.

Solution.

The problem consists of the heat equation (1) and the initial condition (2),

(1) $\frac{\partial u}{\partial t} = c^2 \frac{\partial^2 u}{\partial x^2}$

(2) $u(x, 0) = f(x) \quad (-\infty < x < \infty)$

which in this example is

$$f(x) = U_0 = \text{const} \quad \text{if } |x| < 1 \quad \text{and } 0 \text{ otherwise.}$$

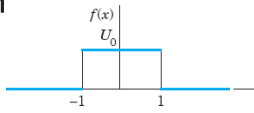


Fig. 298. Initial temperature in Example 1

Our strategy is to take the Fourier transform **with respect to x** and then to solve the resulting **ODE in t**.

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Let $\hat{u} = \mathcal{F}(u)$ denote the Fourier transform of u , regarded as a function of x .

$\frac{\partial u}{\partial t} = c^2 \frac{\partial^2 u}{\partial x^2}$

→

$\mathcal{F}(u_t) = c^2 \mathcal{F}(u_{xx}) = c^2 (-w^2) \mathcal{F}(u) = -c^2 w^2 \hat{u}.$

On the left, assuming that we may interchange the order of differentiation and integration, we have

$$\mathcal{F}(u_t) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} u_t e^{-iwx} dx = \frac{1}{\sqrt{2\pi}} \frac{\partial}{\partial t} \int_{-\infty}^{\infty} u e^{-iwx} dx = \frac{\partial \hat{u}}{\partial t}.$$

Therefore:

$\frac{\partial \hat{u}}{\partial t} = -c^2 w^2 \hat{u}.$

$u(x, t)$

$\hat{u}(w, t)$

Since this equation involves only a derivative with respect to t but none with respect to w , this is a first-order **ODE**, with t as the independent variable and w as a parameter.

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We get the general solution

$$\frac{\partial \hat{u}}{\partial t} = -c^2 w^2 \hat{u} \quad \longrightarrow \quad \hat{u}(w, t) = C(w) e^{-c^2 w^2 t}$$

Initial Condition: $u(x, 0) = f(x)$

$$\hat{u}(w, 0) = C(w) = \hat{f}(w) = \mathcal{F}(f) \quad \longrightarrow \quad \hat{u}(w, t) = \hat{f}(w) e^{-c^2 w^2 t}$$

Now inverse Fourier gives the solution:

$$(14) \quad u(x, t) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \hat{f}(w) e^{-c^2 w^2 t} e^{iwx} dw.$$

$$\hat{f}(w) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(v) e^{-iww} dv.$$

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Combining and assuming that we may invert the order of integration, we then obtain:

$$u(x, t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} f(v) \left[\int_{-\infty}^{\infty} e^{-c^2 w^2 t} e^{i(wx - wv)} dw \right] dv.$$

By the Euler formula

$$e^{-c^2 w^2 t} e^{i(wx - wv)} = e^{-c^2 w^2 t} \cos(wx - wv) + i e^{-c^2 w^2 t} \sin(wx - wv).$$

We see that its imaginary part is an odd function of w , so that its integral is 0.

The real part is an even function of w , so that

$$u(x, t) = \frac{1}{\pi} \int_{-\infty}^{\infty} f(v) \left[\int_0^{\infty} e^{-c^2 w^2 t} \cos(wx - wv) dw \right] dv.$$

This agrees with (9) (with $p = w$) and leads to the further formulas (11) and (13).

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EXAMPLE 3 Solution in Example 1 by the Method of Convolution

Solve the heat problem in Example 1 by the method of convolution.

Solution. The beginning is as in Example 2 and leads to (14), that is,

$$(15) \quad u(x, t) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \hat{f}(w) e^{-c^2 w^2 t} e^{iwx} dw.$$

Since, by the definition of convolution

$$\mathcal{F}(f * g) = \sqrt{2\pi} \mathcal{F}(f) \mathcal{F}(g), \quad (f * g)(x) = \int_{-\infty}^{\infty} \hat{f}(w) \hat{g}(w) e^{iwx} dw,$$

We recognize that this is in the form of **convolution** as:

$$(16) \quad u(x, t) = (f * g)(x) = \int_{-\infty}^{\infty} \hat{f}(w) \hat{g}(w) e^{iwx} dw$$

where

$$(17) \quad \hat{g}(w) = \frac{1}{\sqrt{2\pi}} e^{-c^2 w^2 t}.$$

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by the definition of convolution

$$(18) \quad u(x, t) = (f * g)(x) = \int_{-\infty}^{\infty} f(p) g(x - p) dp,$$

as our next and last step we must determine the inverse Fourier transform g of \hat{g} .

$g(x)$?

Where:

$$\hat{g}(w) = \frac{1}{\sqrt{2\pi}} e^{-c^2 w^2 t}.$$

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$$\hat{g}(w) = \frac{1}{\sqrt{2\pi}} e^{-c^2 w^2 t} \cdot g(x) \quad ?$$

We know (using Fourier Table):

$$\mathcal{F}(e^{-ax^2}) = \frac{1}{\sqrt{2a}} e^{-w^2/(4a)}$$

With $c^2 t = 1/(4a)$ or $a = 1/(4c^2 t)$,

$$\mathcal{F}(e^{-x^2/(4c^2 t)}) = \sqrt{2c^2 t} e^{-c^2 w^2 t} \quad \rightarrow \quad \mathcal{F}(e^{-x^2/(4c^2 t)}) = \sqrt{2c^2 t} \sqrt{2\pi} \hat{g}(w).$$

Hence \hat{g} has the inverse

$$g(t) = \frac{1}{\sqrt{2c^2 t} \sqrt{2\pi}} e^{-x^2/(4c^2 t)}.$$

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Using (18) and replacing x by $x-p$ in $g(t)$ we get:

$$(19) \quad u(x, t) = (f * g)(x) = \frac{1}{2c\sqrt{\pi t}} \int_{-\infty}^{\infty} f(p) \exp\left\{-\frac{(x-p)^2}{4c^2 t}\right\} dp.$$

This solution formula of our problem agrees with (11).

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EXAMPLE 4 Fourier Sine Transform Applied to the Heat Equation

If a laterally insulated bar extends from $x = 0$ to infinity, we can use the Fourier sine transform. We let the initial temperature be $u(x, 0) = f(x)$ and impose the boundary condition $u(0, t) =$

Solution.

$$\frac{\partial u}{\partial t} = c^2 \frac{\partial^2 u}{\partial x^2} \quad \begin{matrix} u(0, t) = 0, \\ u(x, 0) = f(x) \end{matrix}$$

since $f(0) = u(0, 0) = 0,$

(Fourier w.r.t $\rightarrow x$)

$$\mathcal{F}_s(u_t) = \frac{\partial \hat{u}_s}{\partial t} = c^2 \mathcal{F}_s(u_{xx}) = -c^2 w^2 \mathcal{F}_s(u) = -c^2 w^2 \hat{u}_s(w, t).$$

Using:

$$\mathcal{F}_c\{f'(x)\} = w \mathcal{F}_s\{f(x)\} - \sqrt{\frac{2}{\pi}} f(0),$$

$$\mathcal{F}_s\{f'(x)\} = -w \mathcal{F}_c\{f(x)\}.$$

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$$\frac{\partial \hat{u}_s}{\partial t} = -c^2 w^2 \hat{u}_s(w, t).$$

This is a first-order ODE

$$\hat{u}_s(w, t) = C(w) e^{-c^2 w^2 t}.$$

From the initial condition $u(x, 0) = f(x)$ we have $\hat{u}_s(w, 0) = \hat{f}_s(w) = C(w)$. Hence

$$\hat{u}_s(w, t) = \hat{f}_s(w) e^{-c^2 w^2 t}.$$

Taking the inverse Fourier sine transform and substituting

$$\hat{f}_s(w) = \sqrt{\frac{2}{\pi}} \int_0^\infty f(p) \sin wp \, dp$$

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we then obtain the solution formula

(20)
$$u(x, t) = \frac{2}{\pi} \int_0^{\infty} \int_0^{\infty} f(p) \sin wp e^{-c^2 w^2 t} \sin wx dp dw.$$

Figure 300 shows (20) with $c = 1$ for $f(x) = 1$ if $0 \leq x \leq 1$ and 0 otherwise, graphed over the xt -plane for $0 \leq x \leq 2$, $0.01 \leq t \leq 1.5$. Note that the curves of $u(x, t)$ for constant t resemble those in Fig. 299.

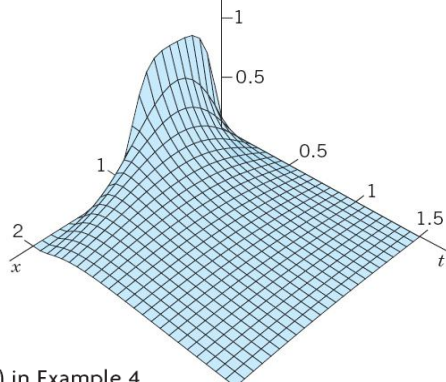


Fig. 300. Solution (20) in Example 4

Questions? Discussion? Suggestions ?



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